

Article

An Assessment of Two Models of Wave Propagation in an Estuary Protected by Breakwaters

Amin Ilia * and James O'Donnell

Department of Marine Sciences, University of Connecticut, Groton, CT 06340, USA; james.odonnell@uconn.edu

* Correspondence: amin.ilia@uconn.edu; Tel.: +1-860-405-9152

Received: 26 July 2018; Accepted: 23 November 2018; Published: 27 November 2018



Abstract: Breakwaters influence coastal wave climate and circulation by blocking and dissipating wave energy. In a large harbor, these effects are combined with wave generation, refraction and reflection. Accurate representation of these processes is essential to the determination of coastal circulation and wave processes. MIKE21SW and SWAN are two third-generation spectral wave models which are used widely in coastal research and engineering applications. Recently improved versions of the models are able to consider the influence of breakwater structures. In this study, we used available observations to evaluate the accuracy of model simulations of waves in New Haven Harbor, Connecticut, USA, an estuary with three detached breakwaters near the mouth. The models were executed on their optimum unstructured triangular grid. The boundary conditions were derived from a bottom mounted Acoustic Doppler Current Profilers (ADCP) on the offshore side of the breakwaters. Wind forcing was applied using data from the Central Long Island Sound buoy. We found that both models were largely consistent with observations during storms. However, MIKE21SW predicted some of storm peaks slightly better. SWAN required the finer grid to achieve the optimum condition, but as it uses a fast, fully implicit algorithm, the computational times were similar. Also, the sensitivity analysis represents that wind forcing and the breakwaters have significant impact on the results.

Keywords: wave hindcast; breakwater; harbor; estuary; SWAN; MIKE21SW; unstructured grid

1. Introduction

Breakwaters are used to protect harbors and shorelines from waves and to limit coastal erosion. In harbors, breakwaters provide tranquility behind them to ease both navigation and berthing for vessels. Breakwaters influence coastal wave climate by breaking, reflecting, and diffracting wave energy. In a large harbor, the fetch may be sufficient for the local generation of waves to be important. Studying these influences in situ poses a challenge. In recent years, spectral wave models have become more widely used and are important in describing coastal wave behavior. However, the performance and precision of the spectral models in real harbors in the presence of the breakwaters has not been well examined.

SWAN, MIKE21SW, and Wavewatch III are three commonly used spectral wave models in coastal and ocean communities. Wavewatch III is mostly used for deep ocean and open sea applications and is not able to include breakwater structures. On the other hand, the recent versions of SWAN and MIKE21SW are equipped to handle breakwaters.

Although parabolic and elliptic mild slope models as well as Boussinesq models may be more appropriate for the simulation of waves in small harbors, where diffraction is the most important phenomenon, in large embayments and harbors wind-wave generation can be important and spectral models are useful. However, some circumstances should meet to use spectral for wave simulation in harbors [1,2].

In recent years, unstructured grid models have become more popular for modeling complex geometries and are an effective alternative to the grid nesting approach, [3,4]. In this study, we employ unstructured SWAN (v41.20) and MIKE21SW (v2017) to simulate the effects of three coastal breakwaters on the wave field in New Haven Harbor, Connecticut, and then assessed their accuracy and efficiency.

There have been previous comparisons of these models. Strauss et al. [5] compared simulations of waves on a narrow continental shelf in Gold Coast, Australia using SWAN and MIKE21SW. SWAN was executed in the fully spectral mode with a structured grid, while MIKE21SW was executed on a directional decoupled parametric mode with an unstructured grid. The results indicated that both models overestimated significant wave height. Moeini and Etemad-Shahidi [6] applied SWAN and MIKE21SW for hindcasting waves in Lake Erie. The study suggested SWAN simulated significant wave height better, while MIKE21SW simulated wave period and direction better. Conversely, Fonseca et al. [7] suggested that MIKE21SW, SWAN, and STWAVE models have similar behavior and precision when examining performance on the Portuguese coast. Hoque et al. [8] evaluated SWAN and MIKE21SW in the Mackenzie Delta in Beaufort Sea, Canada. They concluded that the results of the models were almost identical, but in water shallower than 7 m for wave height between 2 and 3 m, SWAN simulated the significant wave height better, though the peak period from MIKE21SW was more accurate.

In all previous studies, SWAN was executed on a structured grid and MIKE21SW on an unstructured grid. In this study, both SWAN and MIKE21SW were run on their optimum unstructured grids and evaluated inside New Haven Harbor which is separated from Long Island Sound by substantial breakwaters (Figure 1). The model results are compared to data from two bottom mounted Acoustic Doppler Current Profilers (ADCP) in the harbor for all storm events, and to one outside the harbor for storms with northerly winds. The sensitivity of the model results to characteristics of the wind forcing, and the representation of diffraction, reflection, and the breakwaters are discussed. The simulation skill and the efficiency of the models are summarized and possible explanations for their differences were discussed.

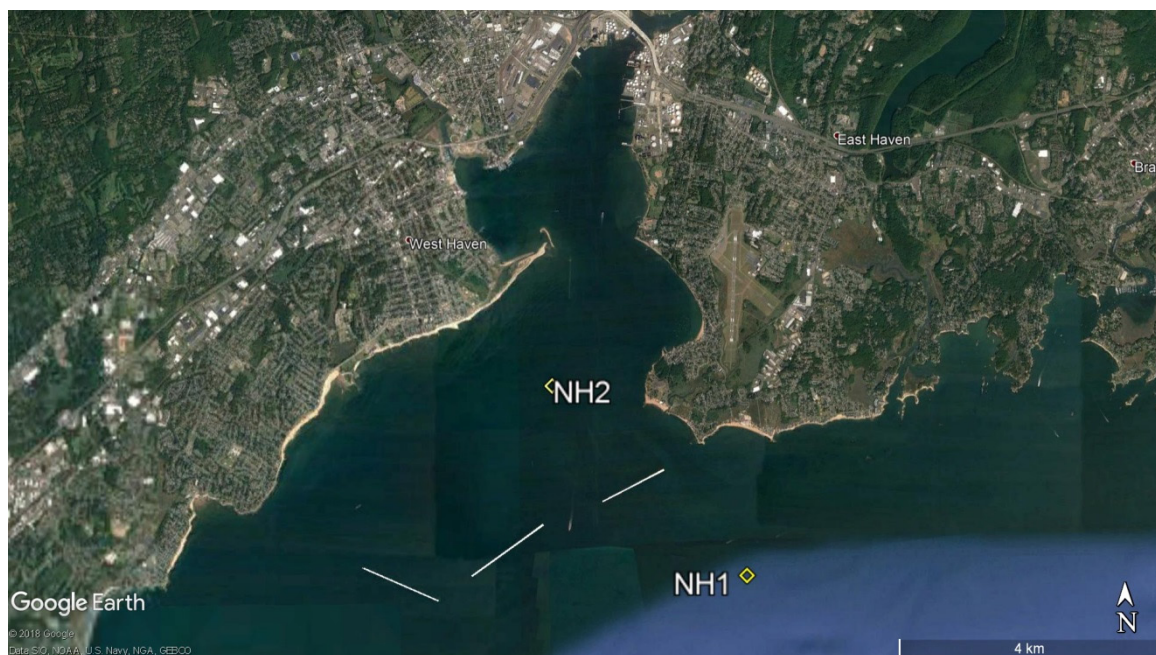


Figure 1. Locations of breakwater structures in front of New Haven Harbor are highlighted by white lines and ADCP locations inside and outside the harbor are identified by yellow boxes and labels. Image Source: “New Haven Harbor”. $41^{\circ} 15'03.79''$ N and $72^{\circ} 55'20.36''$ W. Google Earth. 23 September 2017. 18 June 2018.

2. Model Descriptions

MIKE21SW is a proprietary model developed by DHI Company. It is one of the most widely used wave models in coastal and marine engineering projects around the world. It has a graphical user interface (GUI) that makes it simple to set up the model and visualize the results. SWAN is an open source model developed by the Delft University of Technology. It is often used in academic coastal research and has been integrated to community circulation models. For simplification and clarification, we introduce the models by summarizing their similarities and differences.

SWAN [9] and MIKE21SW [10] are both third generation fully spectral wave models. The models solve the wave action balance equation, described by Mei [11], Komen et al. [12], and Young [13]. In Cartesian coordinates, the wave action balance equation can be written as

$$\frac{\partial N}{\partial t} + \nabla \cdot (\vec{c}_g - U)N = \frac{S}{\sigma} \quad (1)$$

where $N(x, \sigma, \theta, t)$ is the wave action density, $\vec{c}_g = (c_x, c_y, c_\sigma, c_\theta)$ is the wave group velocity, U is ambient current, σ is relative frequency, θ is wave direction, t is time, and S is the total source and sink terms which represent generation, dissipation, and redistribution of wave energy. ∇ is the four-dimensional differential operator with respect to x, y, σ , and θ .

The source terms in both models are almost the same. Wave dissipation terms such as bed friction, wave breaking, whitecapping, nonlinear quadruplet interactions, nonlinear triad wave interactions, and diffraction are essentially the same. However, in some cases such as wind input, whitecapping and quadruplet wave interaction, SWAN provides a wider range of parameterizations and coefficients.

SWAN can be run on both structured and unstructured grids, while MIKE21SW only uses unstructured grids. The spatial discretization method differs between the models. SWAN's spatial discretization is based on a vertex-centered method while MIKE21SW uses a cell-centered method. This implies that wave action N is stored at the grid cell vertices in SWAN and at the cell center in MIKE21SW. Thus, the control volume in SWAN is a polygon, while in MIKE21SW it is triangular.

The numerical methods used in structured and unstructured SWAN are different. In the structured mode of SWAN, a first order upwind-space backward-time (BSBT) scheme, or a second order SORDUP scheme (default for stationary mode), and second order Stelling and Leendertse scheme (default for nonstationary mode) may be selected. In unstructured SWAN the only option is the BSBT scheme which is fast but diffusive [14]. MIKE21SW uses a first order upwind difference and second order accurate scheme in space. The first order scheme is usually sufficient for small-scale domains dominated by local wind. In the case of swell propagation, the second order scheme should be applied [10].

The major difference between the numerical methods is that SWAN uses a fully implicit method time integration whereas MIKE21SW is an explicit approach. Consequently, MIKE21SW avoids solving a large system of equations with the drawback that the temporal step is limited by the Courant number. SWAN's fully implicit scheme eliminates the stability constraint on the temporal step but requires a large system of equations must be solved to achieve a solution. To improve efficiency SWAN uses a point-by-point multi-directional Gauss-Seidel iteration technique that circumvents the need to construct and solve a large system of simultaneous equations as is typical in implicit methods [14]. This technique highly improves the computational efficiency of the SWAN model.

3. Methods

In this study, the model domain is New Haven Harbor, Connecticut, USA. New Haven Harbor is located in Long Island Sound, a large estuary on the northeast coast of the United States. There are three detached breakwaters in front of the harbor to reduce the effects of waves during extreme storms. The area is frequently affected by strong winds during winter, from January to March, and occasionally by hurricanes in the summer and early fall. The significant wave height in central Long Island Sound, where New Haven Harbor is located can exceed two meters. Also, the maximum distance from New

Haven breakwaters to end of estuary is 7 km, sufficient fetch for local wave generation to be important. This is the main reason that the spectral wave models, rather than mild slope or Boussinesq wave models, were employed in this study.

To observe the effect of breakwaters on waves in extreme storms, we deployed two ADCPs, one outside the harbor (NH1) and one inside (NH2), during the winter of 2015 from 21 January to 5 April. The wind data were gathered during the same period by the Central Long Island Sound Buoy (CLIS) which is located 25 km from New Haven. CLIS wind data are a good representative of New Haven wind, particularly when the winds are southerly. Northerly winds may be more influenced by the roughness of the land surrounding the harbor. However, the topography of the New Haven area does not contain features that could cause substantial influence on wind direction. Depths and locations of the observations can be found in Table 1, and Figure 1 illustrates the field site.

Table 1. Location and depth information for the observations used for modeling

Station ID	Latitude (N)	Longitude (W)	Depth (m)	Location
NH1	41°13.44'	72°53.15'	10.4	Outside Harbor
NH2	41°14.64'	72°56.96'	5.2	Inside Harbor
CLIS	41°8.28'	72°39.30'	27	Outside Harbor

Instead of applying a nesting approach, which uses results from a large-scale model to force the boundaries of the local wave model, we used the observed wave spectrum at NH1 to force the open boundary of the models. In this way, we reduce the uncertainty that arises from simulation of the large-scale wave field. The possible variation of the wave field along the boundary was neglected since Long Island Sound is fetch limited and the bathymetric variations near the study site were small. The models were also forced by a uniform wind stress over the domain at half-hour intervals using observation from the CLIS buoy. Mean sea level variations were prescribed in the simulation using observed water level data at the NH2 station.

For the results from the models to be comparable, both models were set up with the same forcing and the optimum time step and grid size for each model were determined by comparison to observations. In MIKE21SW, the user selects a range for time steps (the minimum and maximum) and the model automatically determines the optimum time step based on the grid size, wave propagation speed, and Courant number. The minimum time step must meet the Courant number restriction. In SWAN, the numerical scheme is unconditionally stable but it is critical to ensure that the solution has converged. Using a grid with the resolution range $dx_{max} = 350$ m and $dx_{min} = 35$ m, we found that a time step $dt = 60$ s provided the same solution as a range of smaller time steps, Figure 2a,b compare the solutions at the location of NH2 for two 24 h intervals together with the observations.

We also examined the effect of the grid generation choices. Using a grid converter, available in the MIKE21 package, and the grid generation software in the surface-water modeling system package (SMS v12) we created grids for the models using the same spatial smoothing ratio. Figure 2c,d shows the sensitivity of SWAN solutions for different grid sizes and time steps. As the grid size is decreased, the time step also must be decreased by the same factor to avoid accumulation of numerical computation errors. For the majority of storms, the grid size $dx_{max} = 350$ m and $dx_{min} = 35$ m was adequate (Figure 2c), but some storms required finer grid, (Figure 2d). Refining the grid size and time step by the factor of $\sqrt{2}$ increases the computational time at least by the factor of $(\sqrt{2})^3 = 2.8$; therefore, it is vital that the optimum grid size be selected very carefully. The grid size $dx_{max} = 250$ m and $dx_{min} = 25$ m with time step $dt = 42$ s was found to provide the optimum condition for SWAN.

Figure 2e,f illustrates the sensitivity of the MIKE21SW solutions to grid size. Since MIKE21SW selects the optimum time step based on the grid size automatically, it is easier and faster to get to optimum condition by MIKE21SW. MIKE21SW also showed less sensitivity to grid size than SWAN, (Figure 2e,f). The result was converged for grid size with $dx_{max} = 500$ m and $dx_{min} = 50$ m.

The optimum unstructured triangular grid for each model is presented in Figure 3. The grids were smoothly refined around the breakwaters and inside the harbor.

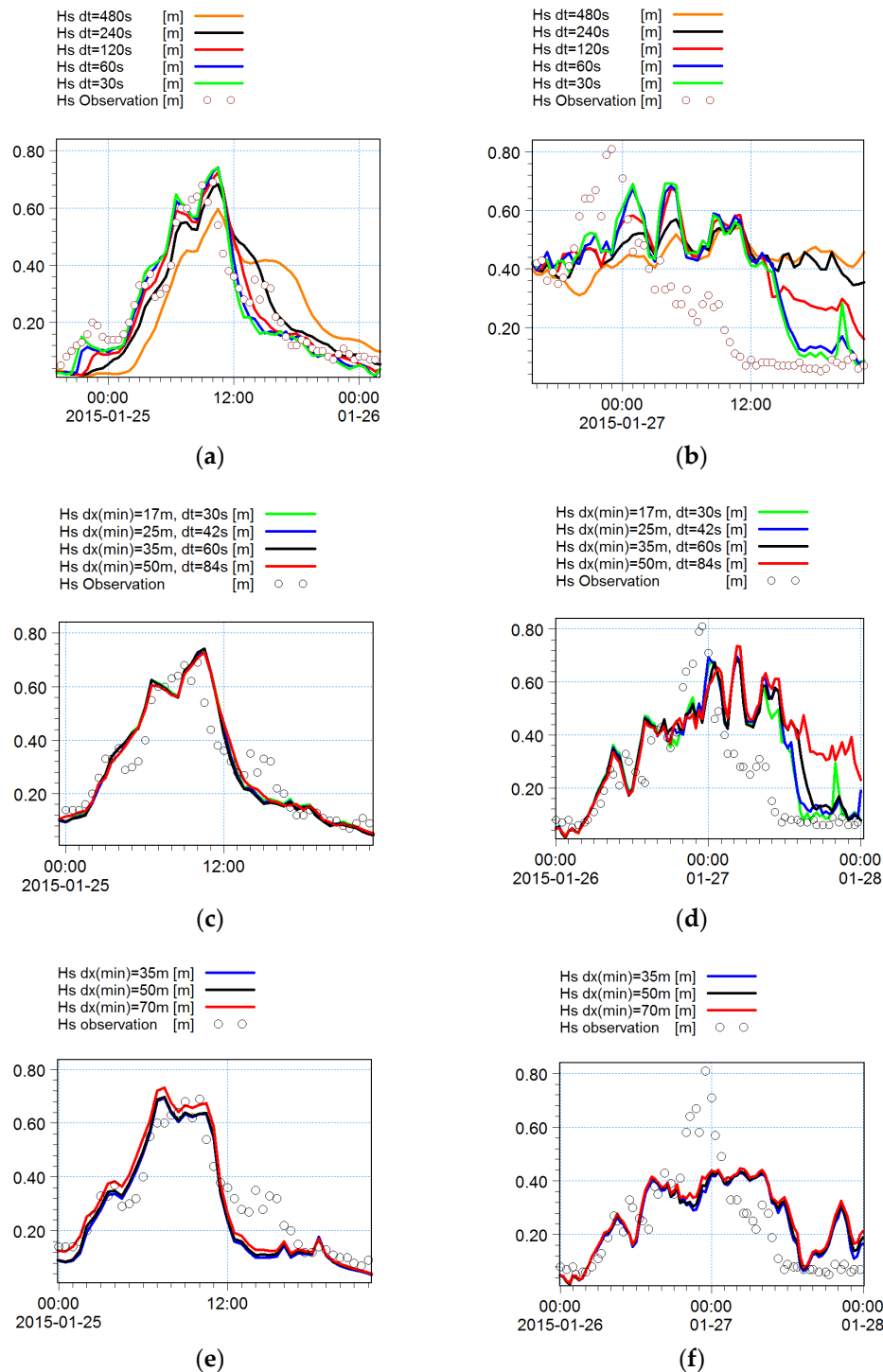


Figure 2. (a,b) Wave height at NH2 from SWAN time step sensitivity analysis using a grid size range $dx_{max} = 350$ m and $dx_{min} = 35$ m. The results indicate that solutions with time steps of 30 s (green) and 60 s (blue) are almost identical for this grid size. (c,d) SWAN solutions for grid size and time step sensitivity analysis, the solution for grid size less than $dx_{max} = 250$ m and $dx_{min} = 25$ m are almost indistinguishable. (e,f) MIKE21SW solutions is less sensitive to grid size and converge with larger grid size $dx_{max} = 500$ m and $dx_{min} = 50$ m.

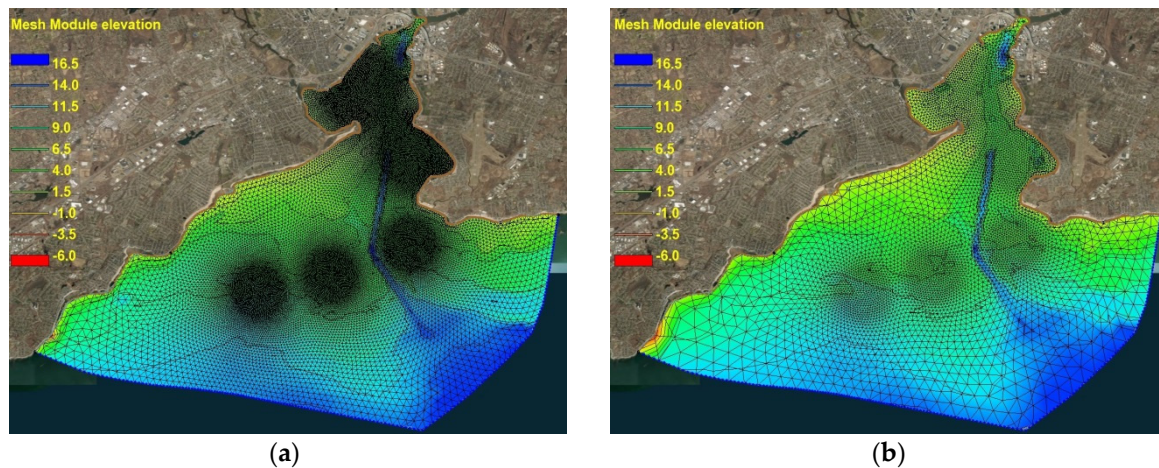


Figure 3. The optimum unstructured triangulate grids used for (a) SWAN with $dx_{max} = 250$ m and $dx_{min} = 25$ m and (b) MIKE21SW with $dx_{max} = 500$ m and $dx_{min} = 50$ m; SWAN required a finer grid to reach the optimum condition. The spatial smoothing of the grids is the same. The grids were refined around the breakwaters, three densification zones, and in the harbor.

The models were calibrated using data from the period of 23 to 28 January. For each model, we tested settings and coefficients, and applied the set that achieved the best agreement with the observations. Both models were executed in third generation fully spectral and non-stationary mode. The same spectral discretization (25 frequencies and 16 directional discretizations) were applied to SWAN and MIKE21SW with a minimum frequency of 0.06 Hz and a maximum of 0.59 Hz.

The results of MIKE21SW were tested for both first and second order numerical scheme, and no significant difference was observed between the first order and second order schemes. Consequently, in order to save computational time, MIKE21SW was run on first order. As mentioned in Section 2, the only available numerical scheme in unstructured SWAN is first order BSBT.

Options for representing the wind input source and whitecapping dissipation functions are different in the models. SWAN is configured to use the wind input source function of Janssen [15,16], as implemented in WAM cycle 4 [17]; the method of Komen [18], as implemented in WAM cycle 3 [19]; and method of Yan [20]. MIKE21SW only supports Janssen's method. Dissipation through whitecapping is based on the development of Hasselmann [21] in both Janssen's and Komen's methods, however, the coefficients are different (see Komen [12,18]). Yan's wind input method is combined with saturation-based whitecapping as described in Van der Westhuysen [22,23]. Moeini and Etemad-Shahidi [6] suggested that Komen's method led to more accurate significant wave height than Janssen's method for Lake Erie. Hoque [8] indicated Westhuysen's formulation tends to have better significant wave height in the Mackenzie Delta. We tested all three methods with SWAN. We found that Janssen's method provides the best wave height simulations in the harbor. Therefore, Janssen's wind input method was used for both models with the tunable coefficients set as $C_{ds}^* = 4.5$ and $\delta = 0.5$.

Wave dissipation due to bottom friction was represented in both models by the empirical JONSWAP [24] approach. The friction coefficients introduced by Zijlema et al. [25] were used: $C_b = 0.038 \text{ m}^2\text{s}^{-3}$ for SWAN and $c_{fw} = 0.0077 \text{ ms}^{-1}$ for MIKE21SW.

Waves breaking in shallow water was taken into account in the models using the Battjes and Janssen [26] formulation with $\alpha = 1$ and $\gamma = 0.8$. Also, triad wave-wave interactions were enabled in both models. Triad interactions in MIKE21SW are calculated based on Eldeberky and Battjes [27] and in SWAN based on Eldeberky [28], a slightly modified version of the former. Quadruplet wave-wave interactions were also enabled in the models. In both models, the quadruplet wave-wave interactions are computed with the Discrete Interaction Approximation (DIA) as proposed by Hasselmann et al. [29].

When breakwaters are present in the model domain, diffraction computations become important. Diffraction is taken into account in the models using a phase-decoupled refraction-diffraction approximation proposed by Holthuijsen et al. [1]. Diffraction computation in the models is almost the same, but a wave field smoothing technique for the computation of the diffraction parameter that is not available in SWAN with an unstructured grid.

The reflection coefficients for breakwaters were calculated using the method proposed in the Coastal Protection Manual: Part VI, [30], and originally developed by Seeling [31], for non-overtopped slopping structures with the parameter values of Davidson [32]. The reflection coefficient for the type of structure and wave climate in the study area varied from 0.47 to 0.52. Therefore, the average reflection coefficient of 0.493 was selected.

Spectral wave models adequately simulate the effects of diffraction when breakwaters are far from the coastline and have a low reflection coefficient [1]. In addition, the breakwaters should not cover the down-wave view substantially, [1]. In New Haven harbor, the distance between the breakwaters and coastline varies 4 to 5.2 km, and the separation of the breakwaters is 700 m to 1000 m, which much larger than the threshold distance [2] of twice the dominant wavelength (here $2L \cong 100$ m, where L is wave length). Also, the rubble mound breakwaters with low reflection coefficient, 0.493, create incoherent wave reflection.

The spectral wave models such as SWAN simulate wave diffraction better for the wider directional spectrum of wind-waves than for swell [2]. The wave spectra for the important storms are shown in Figure 4. The wave spectra in the harbor are directionally broad, similar to those obtained in Long Island Sound. Therefore, New Haven harbor is considered a suitable case for simulating waves using the spectral models.

4. Results

The results of SWAN and MIKE21SW models were compared and assessed inside the harbor at NH2 during all storms and at NH1 for northerly storms. We divided the observations into five storm periods, detailed in Table 2. Each of these periods is discussed separately, and then the results are summarized. The statistical parameters used for data validation are

$$Bias = \frac{\sum_{i=1}^n (Y_i - X_i)}{n}, \quad (2)$$

$$RMS = \sqrt{\frac{\sum_{i=1}^n (Y_i - X_i)^2}{n}}, \quad (3)$$

$$m = \text{slope of the best fitted line}, \quad (4)$$

$$R^2 = 1 - \frac{\sum_{i=1}^n (Y_i - Y'_i)^2}{\sum_{i=1}^n (Y_i - \bar{Y})^2}, \quad (5)$$

where Y_i , with mean \bar{Y} , are the observed values, X_i are the simulated values, and n is the number of data points. RMS is the root mean square error, R^2 is the fraction of the variance in the data explained by the model, Y'_i is the estimated value by regression, and \bar{Y} is the mean of the observed values.

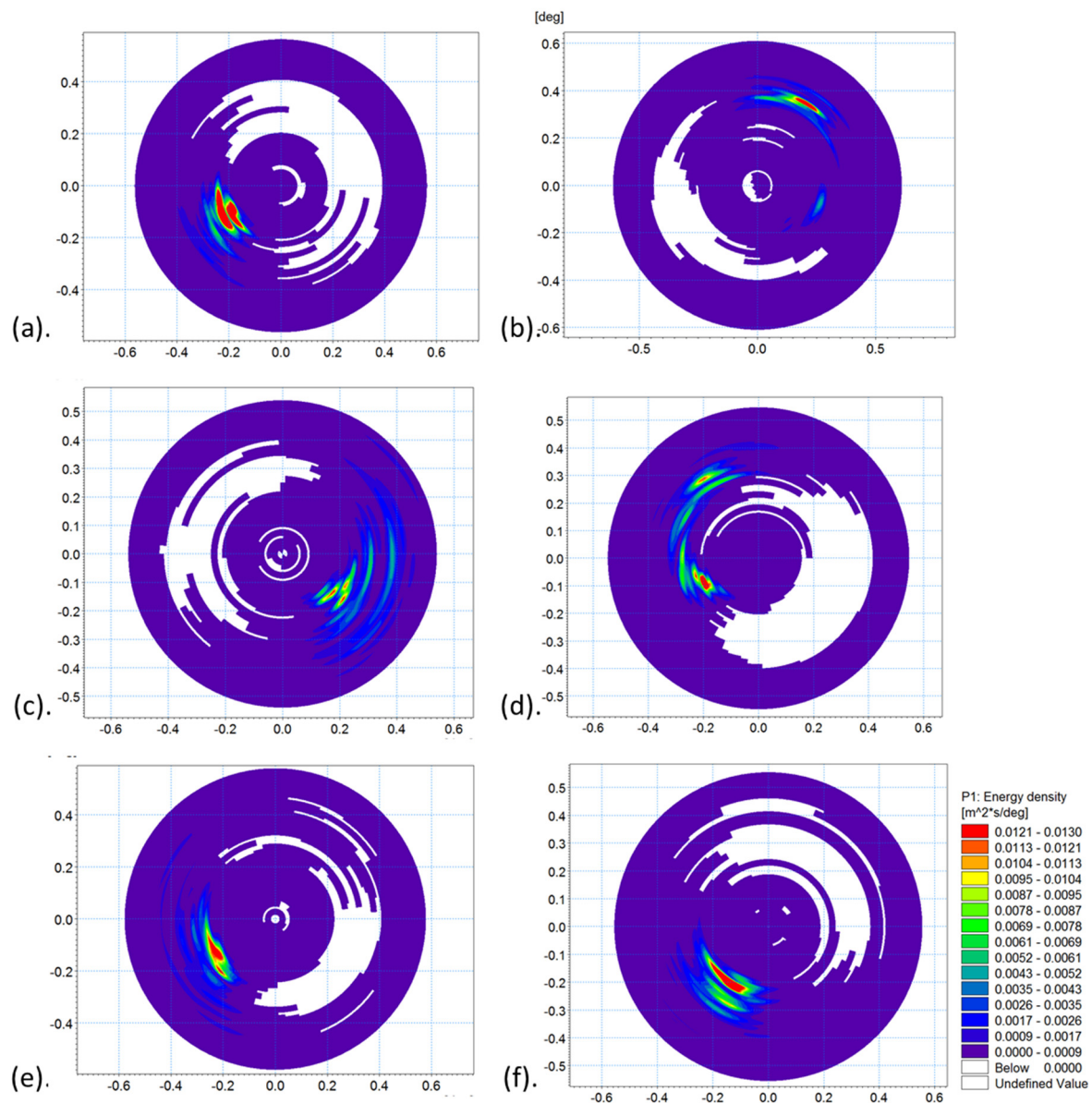


Figure 4. The two-dimensional wave spectra for important storms at NH1. As the majority of waves generates with local wind in central Long Island Sound, the spectrums are directionally wide. (a) 25 January; (b) 27 January; (c) 2 February; (d) 15 February; (e) 7 March; and (f) 2 April. The spectral wave models simulate wave diffraction for directionally wide spectrum better than directionally narrow spectrum, [2].

Table 2. Storm periods used for analyzing and assessing the models

No.	Storm Period
1	23 to 28 January
2	1 to 10 February
3	14 to 21 February
4	6 to 8 March
5	2 to 5 April

Three consecutive high wind intervals occurred during the first storm period, 23 to 28 January. Wave boundary condition and wind input data, as well as comparisons between observations and the model predictions of the significant wave height, peak wave period and wave direction for Storm

Period 1 are presented in Figure 5. The noise in observed wave periods and directions in low significant wave heights events were not included in the assessments. In this storm period, the first two storms had winds from the west and southwest. In the first storm, the predicted significant wave height (H_s) series were very similar to each other. Both missed the first part of the storm peak and under-predicted the second part by 5% (Figure 5b). Both models were able to simulate wave period and direction correctly (Figure 5c,d). In the second storm, both models reproduced the storm's entire wave height peak very well. Simulated wave period and direction for the second storm also agreed well with the observations. For the third storm, the models did not perform well. On 26 January, the storm caused easterly winds, then on 27 January, wind stress vectors suddenly turned by 90 degrees and a very strong wind (in excess of 15 m/s) blew from the north. This strong wind was adequate to generate wave in a very fetch limited area as New Haven estuary. The models showed different behavior during this storm. Simulated peak wave periods were in consistent with observations. Modeled wave directions for the first part of the storm completely agreed with observation, but after the wind vector rotated on 26 January 10 p.m., there was a 25 to 40 degrees difference between observed and simulated wave direction. The observation showed a wave direction from northeast (50 degrees), while simulated wave directions were from north. The statistical parameters show that SWAN results were slightly better than MIKE21SW for this period, (Table 3). The *Bias* in significant wave heights obtained from SWAN was 0.03 m less than that in MIKE21SW. R^2 for SWAN was 0.60 versus 0.57 for MIKE21SW. The *RMS* values were the same (0.13) and the slope of best fitted line was better for MIKE21SW (0.89) than SWAN (0.72).

Comparison of the results of the models and observations from Storm Period 2, 1 to 10 February, are shown in Figure 6. Both models showed similar behavior during this period though SWAN overestimated the strongest storm significant wave height on 2 February by 50% (Figure 6b). There was a high wind event from the north on 5 and 6 February (Figure 6a), and the models simulated the significant wave height very well (Figure 6b). In other storms, the models had very similar behavior, both missed some small oscillations and overestimated others. The models correctly simulated the peak wave period and mean wave direction, (Figure 6c,d). Statistical parameters for this period did not present any substantial preference of the models, *Bias* and R^2 were better for SWAN and *RMS* and *m* were better for MIKE21SW (Table 3).

Three storms occurred in the third period, from 14 to 21 February. As shown in Figure 7, the models accurately simulated the first storm. The second storm started with a wind from north to south and it then rotated to the east. MIKE21SW accurately simulated significant wave height in the first part of the storm but later produced overestimates. In contrast, SWAN overestimated significant wave height over the whole storm duration. Both models provide poor estimates of the wave period for the second part of the storm. For the third storm, both models successfully simulated the wave height, period, and direction, however, MIKE21SW slightly under-predicted the first part. Overall, *RMS* error was 0.03 lower, the slope of the best-fitted line was 0.11 higher, and R^2 was 0.01 higher for MIKE21SW. The magnitude of *Bias* values were the same (0.02) but with opposite signs, (Table 3). Therefore, MIKE21SW results were slightly more accurate than SWAN for this time period, mostly due to better simulation of the second storm.

In the fourth storm period, both models failed to correctly simulate significant wave height variations between 6 and 8 March, as shown in Figure 8. The storm includes two peaks in significant wave height, but the models underestimated both (Figure 8b). However, they correctly estimated wave periods and directions (Figure 8c) and the statistical parameters show similar model performance. There was two hours delay in storm growth in both models. The storm started at 3 March 03:00 p.m., according to observations, but 05:00 p.m. in the models. Also, there was no difference between wind speed at 03:00 and 05:00 p.m. (both around 5 m/s, Figure 8a). Some statistical parameters, such as R^2 , were better for MIKE21SW and others, such as *Bias*, for SWAN (Table 3).

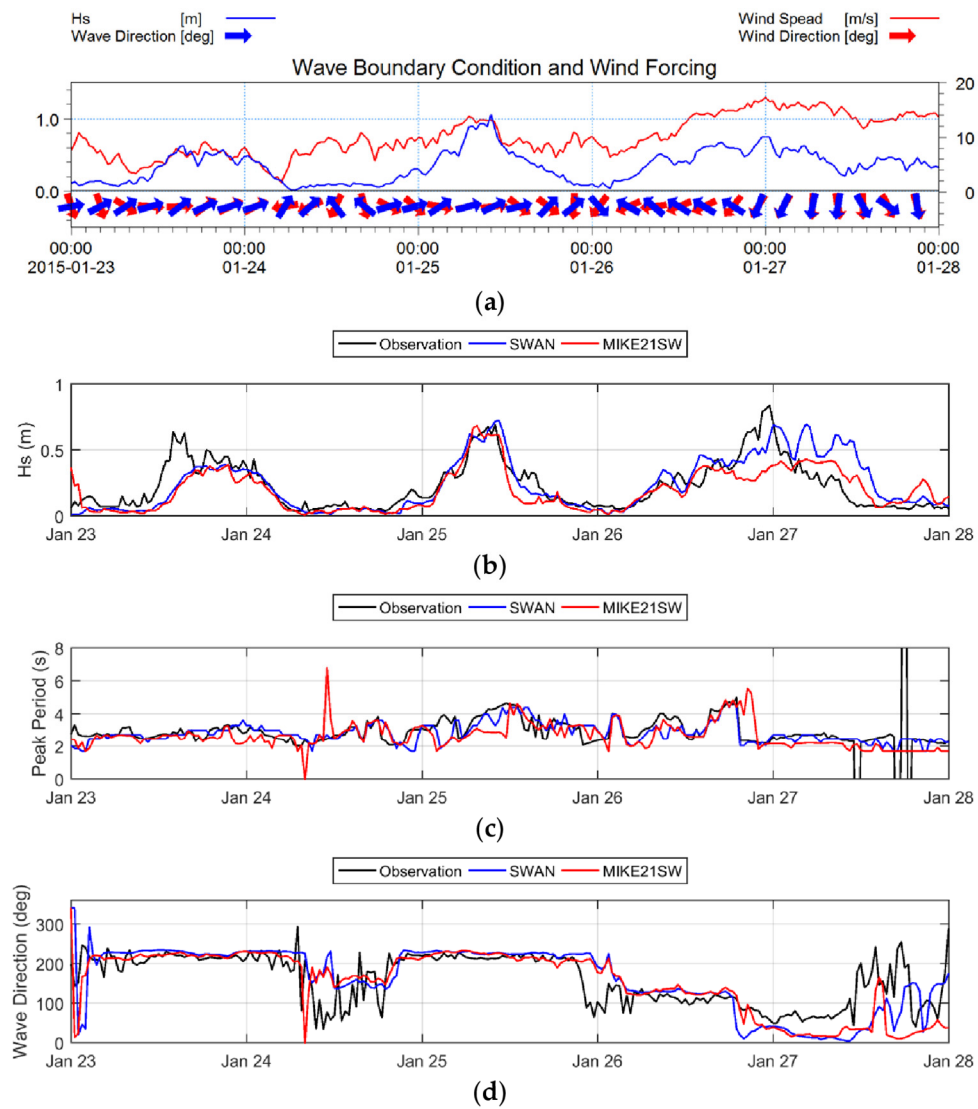


Figure 5. (a) Wave boundary condition and wind input for the time interval of 23 to 28 January. Blue: significant wave height (left axis), blue arrow: wave direction at the boundary, red: wind speed (right axis), red arrow: wind direction. (b–d) Comparison of SWAN (blue), MIKE21SW (red), and observation (black) at NH2 for the time interval of 23 to 28 January. (b) Significant wave height. (c) Peak wave period (d) Mean wave direction. The models had good performance for the first two southerly events but not for the third event which was from the north.

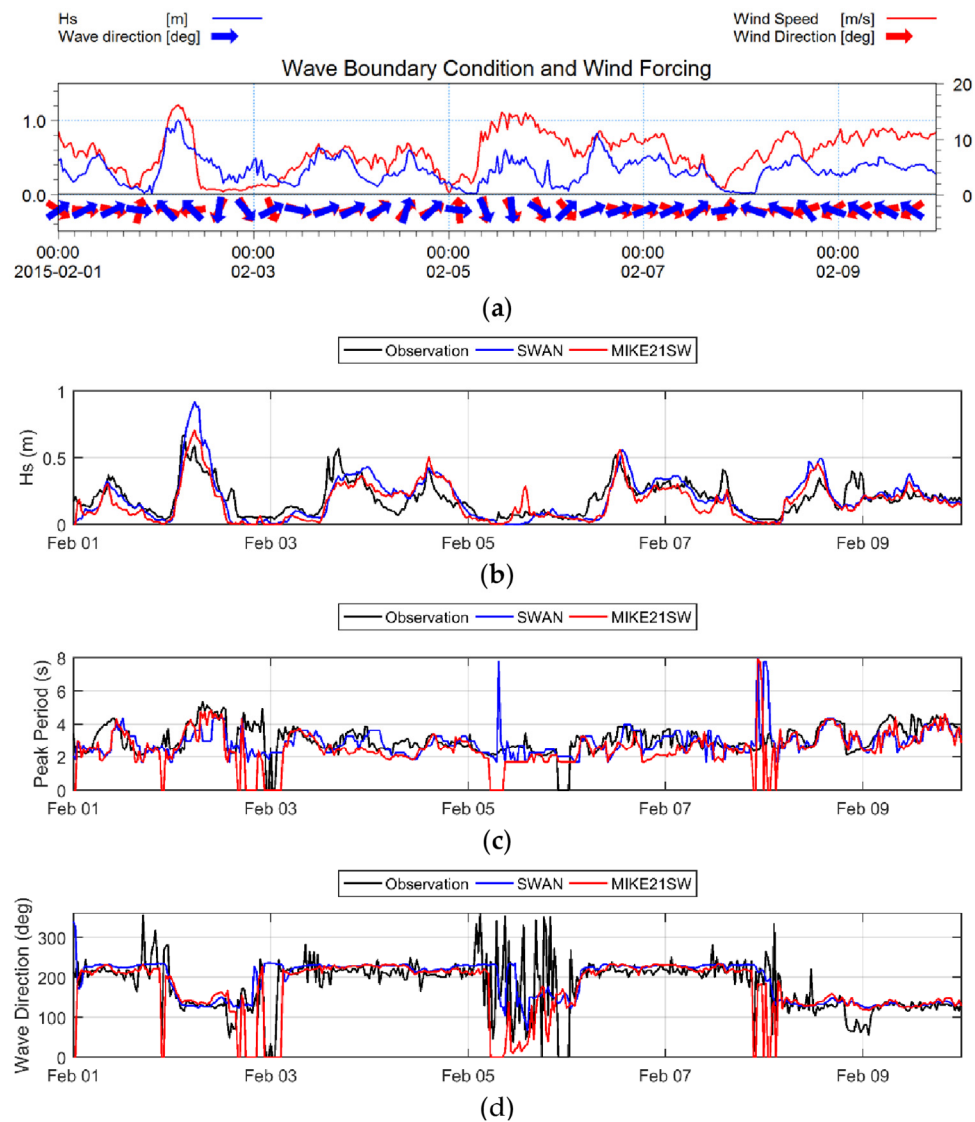


Figure 6. (a) Wave boundary condition and wind input for the time interval of 1 to 10 February. Blue: significant wave height (left axis), blue arrow: wave direction at the boundary, red: wind speed (right axis), red arrow: wind direction. (b–d) Comparison of SWAN (blue), MIKE21SW (red), and observation (black) at NH2 for the time interval of 1 to 10 February. (b) Significant wave height. (c) Peak wave period (d) Mean wave direction. The models had good performance for this period though SWAN overestimated the strongest storm on 2 February.

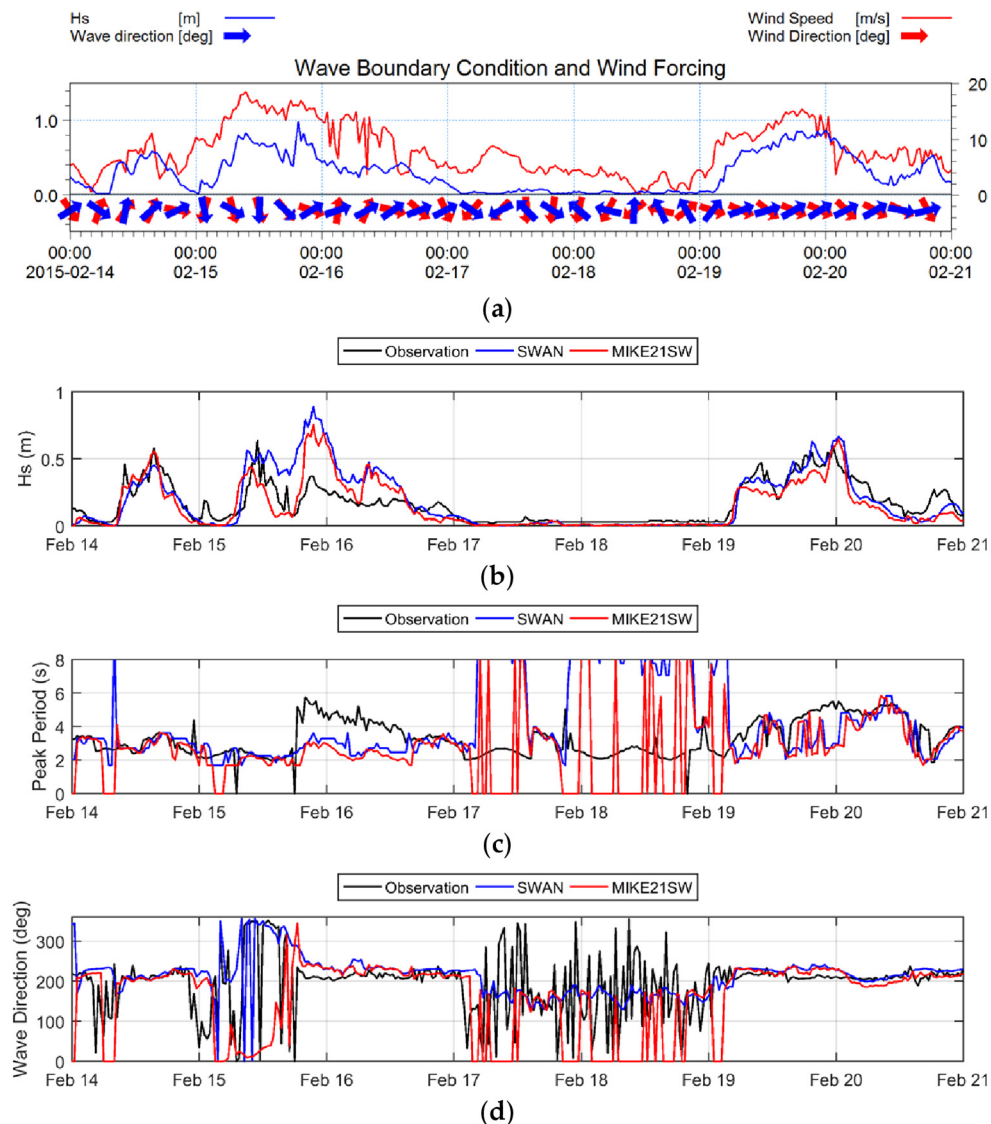


Figure 7. (a) Wave boundary condition and wind input for the time interval of 14 to 21 February. Blue: significant wave height (left axis), blue arrow: wave direction at the boundary, red: wind speed (right axis), red arrow: wind direction. (b–d) Comparison of SWAN (blue), MIKE21SW (red), and observation (black) at NH2 for the time interval of 14 to 21 February. (b) Significant wave height. (c) Peak wave period (d) Mean wave direction. During the second storm, when the wind direction changed from north to south to west to east, the simulations were poor.

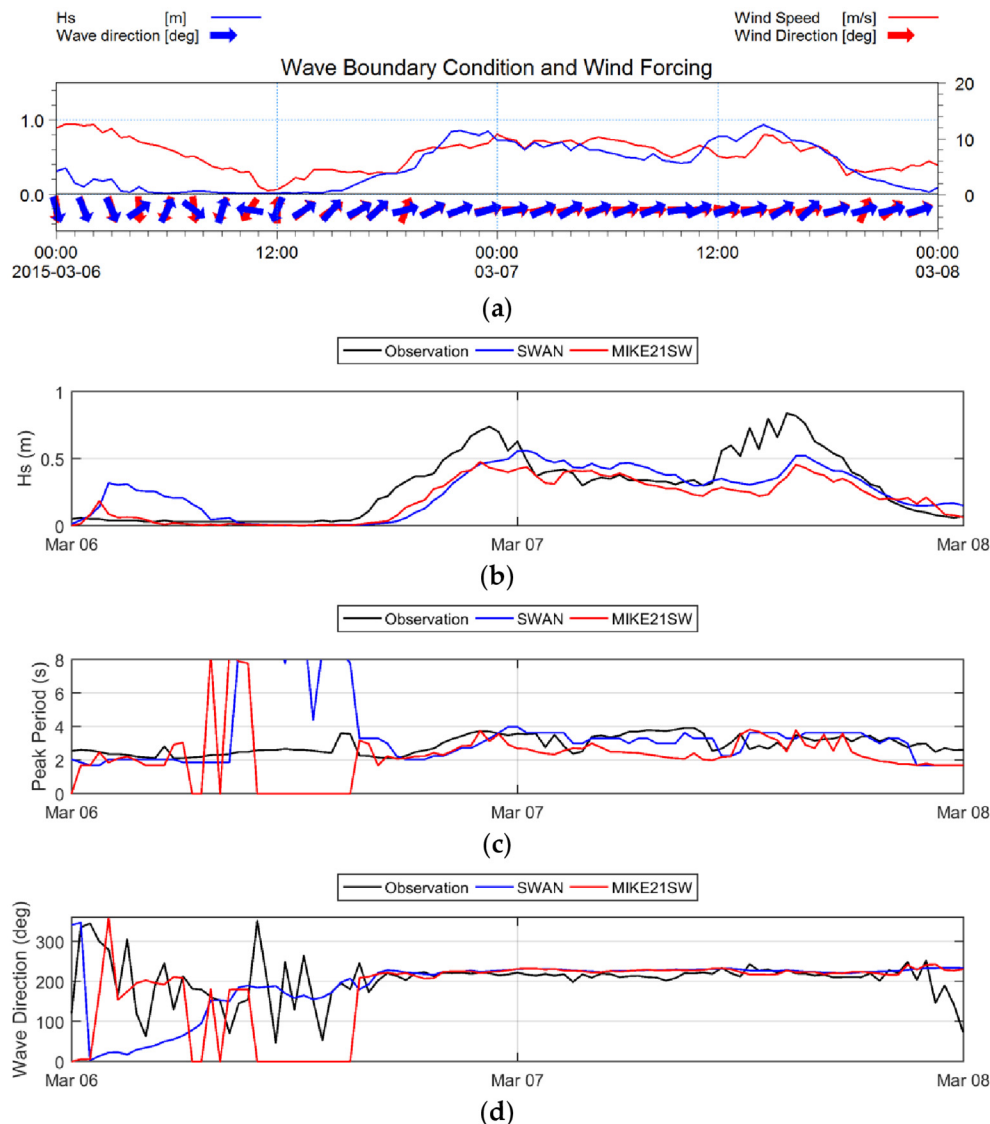


Figure 8. (a) Wave boundary condition and wind input for the time interval of 6 to 8 March. Blue: significant wave height (left axis), blue arrow: wave direction at the boundary, red: wind speed (right axis), red arrow: wind direction. (b–d) Comparison of SWAN (blue), MIKE21SW (red), and observation (black) at NH2 for the time interval of 6 to 8 March. (b) Significant wave height. (c) Peak wave period (d) Mean wave direction. Both models did not do well during this time interval. A two-hour delay in storm growth was significant in models results.

Results from Storm Period 5 are illustrated in Figure 9, showing the assessment of the models results during the storm took place on 2 and 3 April. MIKE21SW was able to catch the storms highest wave height and correctly computed wave period and direction during the storm. SWAN slightly underestimated the first part of the storm and there was a short delay in storm growth (Figure 9b). Statistical parameters were slightly better for MIKE21SW with $R^2 = 0.81$ versus SWAN results with $R^2 = 0.77$. RMS was better for MIKE21SW and the slope of best-fitted line was better for SWAN (Table 3).

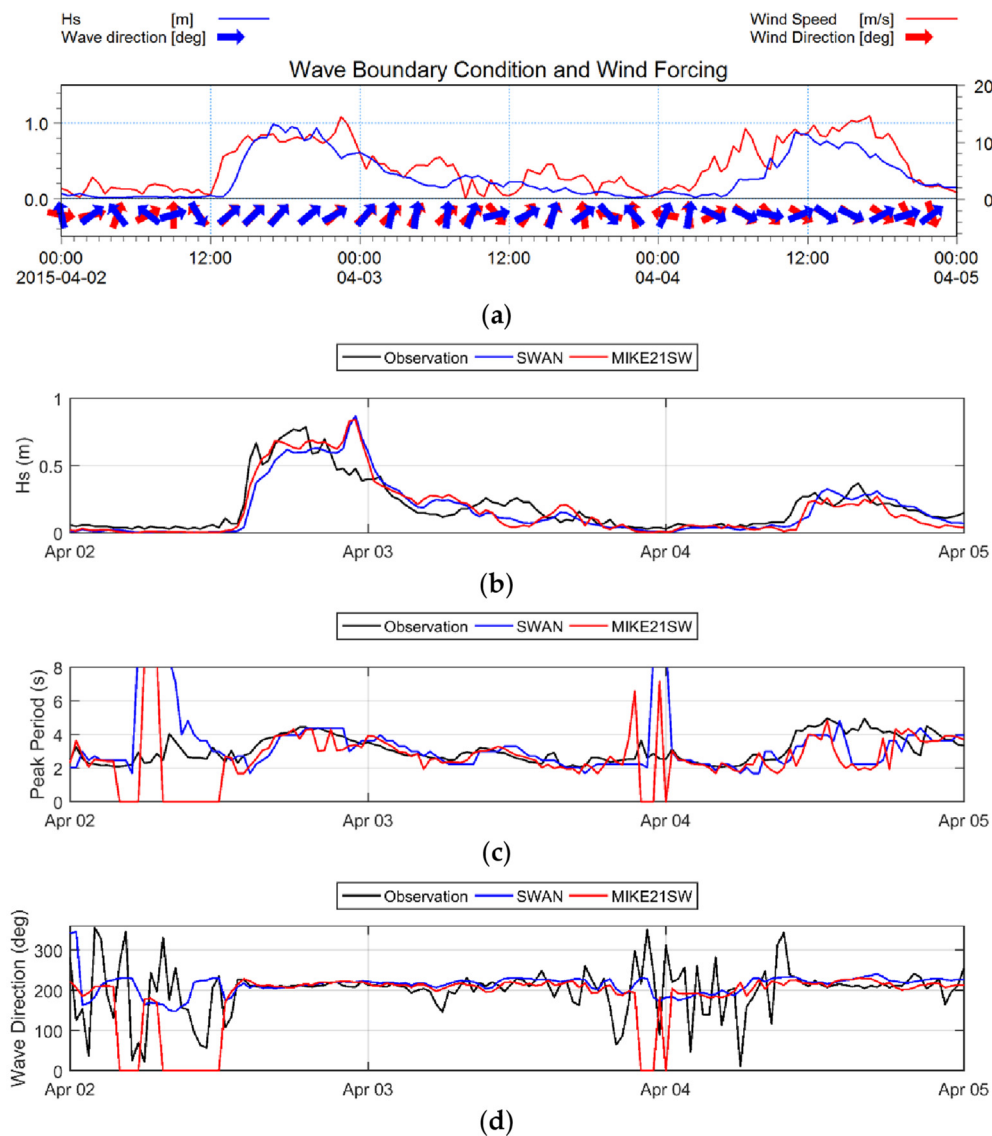


Figure 9. (a) Wave boundary condition and wind input for the time interval of 2 to 5 April. Blue: significant wave height (left axis), blue arrow: wave direction at the boundary, red: wind speed (right axis), red arrow: wind direction. (b–d) Comparison of SWAN (blue), MIKE21SW (red), and observation (black) at NH2 for the time interval of 2 to 5 April. (b) Significant wave height. (c) Peak wave period (d) Mean wave direction. Both models performed well for this storm period, but MIKE21SW was slightly better during the first part of the storm on 2 March.

We compare all storm simulations to data in Figure 10. *Bias* was better for SWAN. *RMS* of MIKE21SW was 0.01 lower than *RMS* obtained from SWAN results. The slope of best-fitted line for MIKE21SW was closer to one. R^2 of the models' results were almost the same, 0.61 for MIKE21SW versus 0.60 for SWAN. Therefore, the statistical parameters suggest the performance of both models were good and the models' results had a lot of similarity.

The model predictions did not agree with observation at NH2, inside the harbor, for two northerly storms, 27 January and 15 February. To assess whether the errors were due to the wind speed magnitude being too high, we compared the results of the models to data obtained at NH1, at the boundary of the model domain, (Figure 11). Note that since wave energy was propagating out of the domain during northerly wind the boundary observations were not influencing in the predictions. Both model results were biased high relative the observation at NH1 with the SWAN results having a larger bias (Figure 11a). However, both models were more correlated with observations at NH1

than at NH2. Both models did a better job for the second northerly event on 15 February. SWAN showed lower errors early in the simulation but both models overestimated the significant wave height later in the storm just as they did in the interior of the Harbor at NH2. This may be a consequence of overestimation of the wind stress, however, the reduced correlation in the model and data time series cannot be explained by the magnitude bias alone.

Table 3. Comparison of statically parameter between SWAN and MIKE21SW. Statistically, there was no substantial difference between the performances of the models. SWAN performance was better for the storm periods 23 to 28 January while MIKE21SW for 14 to 21 February and 2 to 5 April.

Storm Period	SWAN					MIKE21SW				
	<i>n</i>	<i>Bias</i>	<i>Rms</i>	<i>m</i>	R^2	<i>n</i>	<i>Bias</i>	<i>Rms</i>	<i>m</i>	R^2
23 to 28 January	241	−0.01	0.13	0.72	0.60	241	0.04	0.13	0.89	0.57
1 to 10 February	433	0.00	0.10	0.60	0.62	433	0.03	0.09	0.69	0.60
14 to 21 February	337	−0.02	0.14	0.55	0.60	337	0.02	0.11	0.66	0.61
6 to 8 March	97	0.03	0.16	1.03	0.56	97	0.09	0.16	1.33	0.73
2 to 5 April	145	0.02	0.10	0.83	0.77	145	0.02	0.09	0.79	0.81
All Storms	1253	0.00	0.12	0.67	0.60	1253	0.03	0.11	0.79	0.61

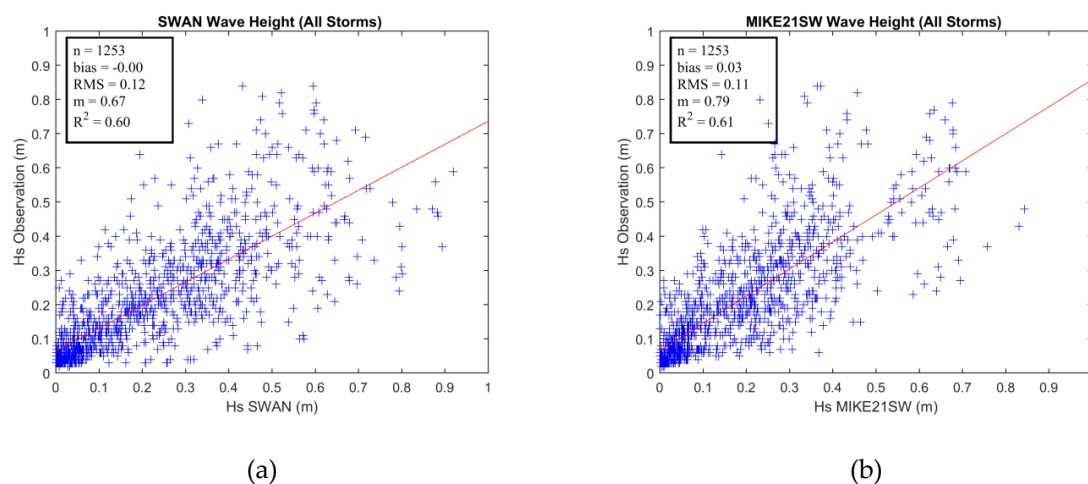


Figure 10. Comparison of scatter plot and statically parameter between observed and simulated significant wave heights from SWAN (a) and MIKE21SW (b) for all storms. There was no substantial difference between the performance of the models. *Bias* was better for SWAN and *RMS*, the slope best fitted line and R^2 were better for MIKE21SW.

Figure 12 displays the variation of significant wave height over the model domain for three storms on 25 January (wind from the southwest), 27 January (wind from the north), and 2 February (wind from the southeast). Unfortunately, as the output format of the models were different, we were not able to plot them with the same tools, therefore, the color bars scales are slightly different. The models had similar behavior over the domain on the southern storms (Figure 12a,b), but different on the northern storm (Figure 12c,d). Also, it can be implied that MIKE21SW dissipates wave energy around the breakwaters more than SWAN, it may be the reason or the SWAN overestimation on 2 February.

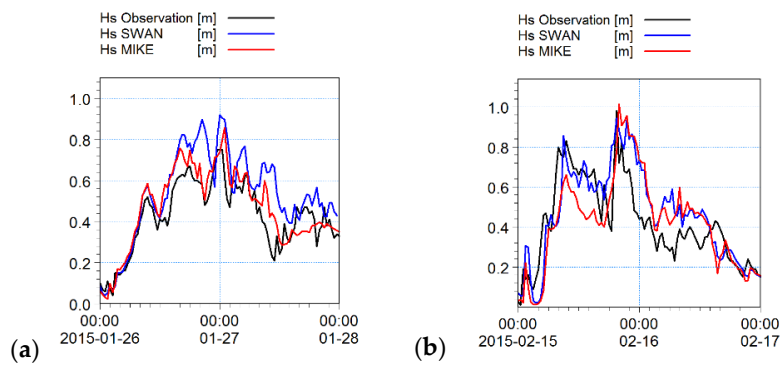


Figure 11. Comparison of significant wave height for two northerly storms at NH1 station outside the harbor. SWAN (blue), MIKE21SW (red), and observation (black) (a) 26 to 28 January. (b) 15 to 17 February.

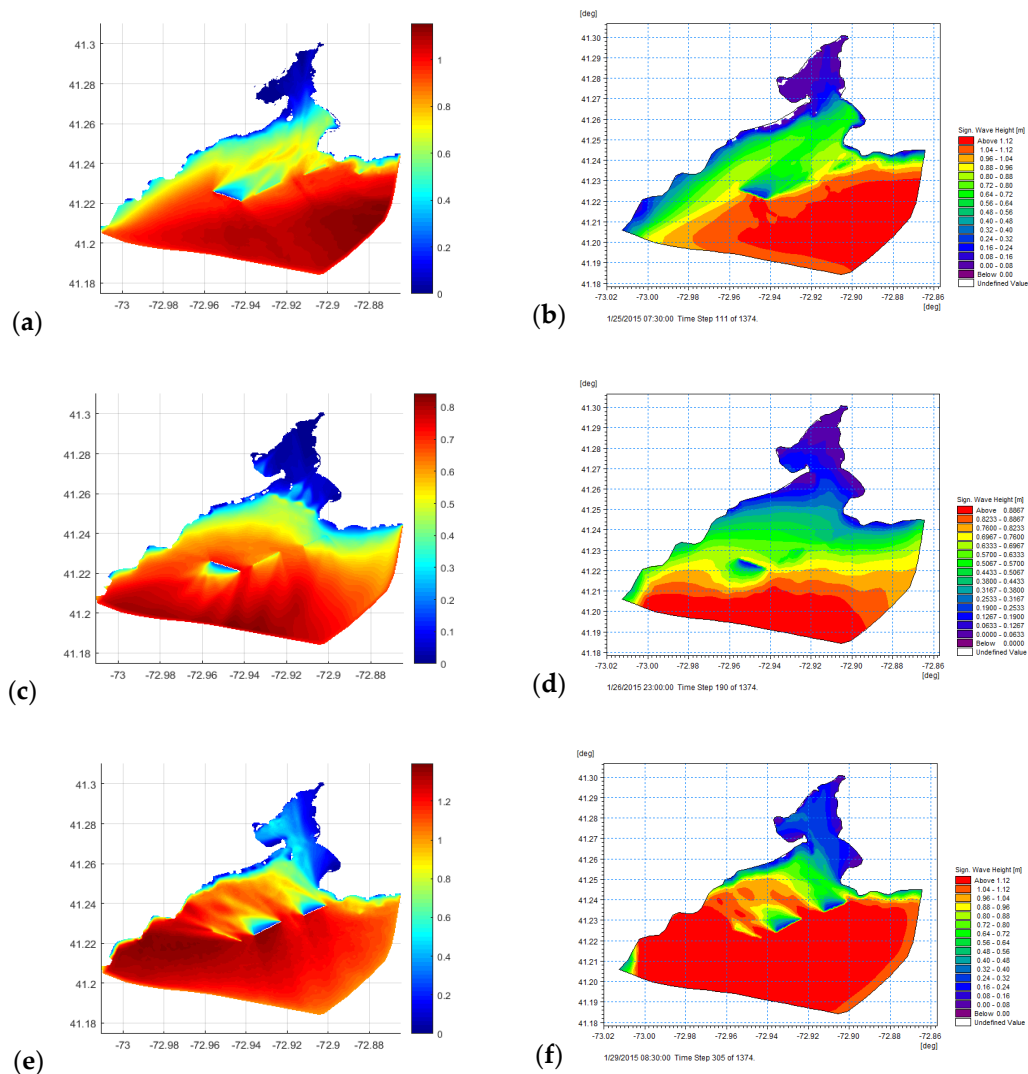


Figure 12. The variation of significant wave height over modeling domain for three main storms from the southwest, north, and southeast. (a) SWAN, 25 January, storm from the southwest; (b) MIKE21SW, 25 January, storm from the southwest; (c) SWAN, 27 January, storm from the north; (d) MIKE21SW, 27 January, storm from the north; (e) SWAN, 2 February, storm from the southeast; (f) MIKE21SW, 2 February, storm from the southeast. Note: the color bar scales are slightly different for SWAN and MIKE21SW, as they were plotted by different tools.

5. Discussion

In total, 14 storms occurred in New Haven harbor during the observation period in the winter of 2015 and the models simulated similar wave fields in most of them. Both models had very good statistical performance when storm winds blew from the south, when the breakwaters are most influential. SWAN predicted the peak significant wave height for one storm (on 18 February) better than MIKE21SW, while MIKE21SW simulated three storms (on 2 February, 15 February, and 3 April) better than SWAN. The worst performance during southerly wind storms was the SWAN results for 2 February when the model overestimated the peak significant wave height by 50%. Figure 12e suggest that the wave sensor was close to a region of high spatial gradient in the significant wave height during this storm and, therefore, slight differences in diffraction and propagation led to large differences in model solution values. The worst performance of MIKE21SW was on 18 February when the peak significant wave height was underestimated by 25%. In other storms the models had similar behavior. Notably, both models underestimated the storm peak on 7 March by 60%.

To further understand the sensitivity of the models results to different physical processes we ran the simulations with the wind forcing eliminated; zero reflection from the breakwaters; with diffraction disabled; and with the breakwaters entirely removed. Figure 13 shows the results of these simulations at NH2.

Figure 13a,d show the observations (+symbols) and solutions at NH2 for the 25 January storm with SWAN and MIKE21SW when the wind blew from the south. The agreement between the black lines and the +symbols shows that both models performed well. The differences between the black and blue lines shows that the influence of the wind over the harbor is significant and increases the peak significant wave height during the storm by 50%. Comparison of the black line to the magenta, green, and red lines show that the next most important process is the presence of the breakwaters. Eliminating them increases the peak significant wave height by approximately 15% in SWAN and 30% in MIKE21SW. Figure 13b,e shows the same properties for the 2 February storm period when the wind was from the south east. Since the performance of both models were not as good, due to high spatial gradients near the location of the wave sensor, the comparison of the result with (black line) and without local wind forcing (blue line) shows again that the local wind can increase the peak significant wave height by approximately 50%. Similarly, removing the breakwaters increases the peak wave heights as in Figure 13a,d, by comparable amounts. We note that the presence of breakwaters appears to be more significant in MIKE21SW than SWAN. The spatial distributions of wave height in Figure 12 also indicate this difference in model performance.

Comparison of the green and red lines with the black lines in Figure 13a,b,d,e illustrates that the effects of diffraction and reflection influenced the significant wave height at NH2 by less than 5 percent in both models. NH2 is approximately 2 km (or $40L$, where L is the dominant wave length) away from the breakwaters. This is consistent with [2] which concluded that reflection and diffraction effects are insignificant far away the breakwaters.

Figure 13c,f show the data and solutions for the 27 January storm when the wind was from the north. As expected, comparison of the black and blue lines shows that in the absence of wind forcing the significant wave height drops to zero. The other effects do not play a large role in the model predictions though reflection, also, has a maximum effect of 7%. These results highlight the importance of wind in this case study, which is a large harbor with fetch length varies from 5 km to 7 km.

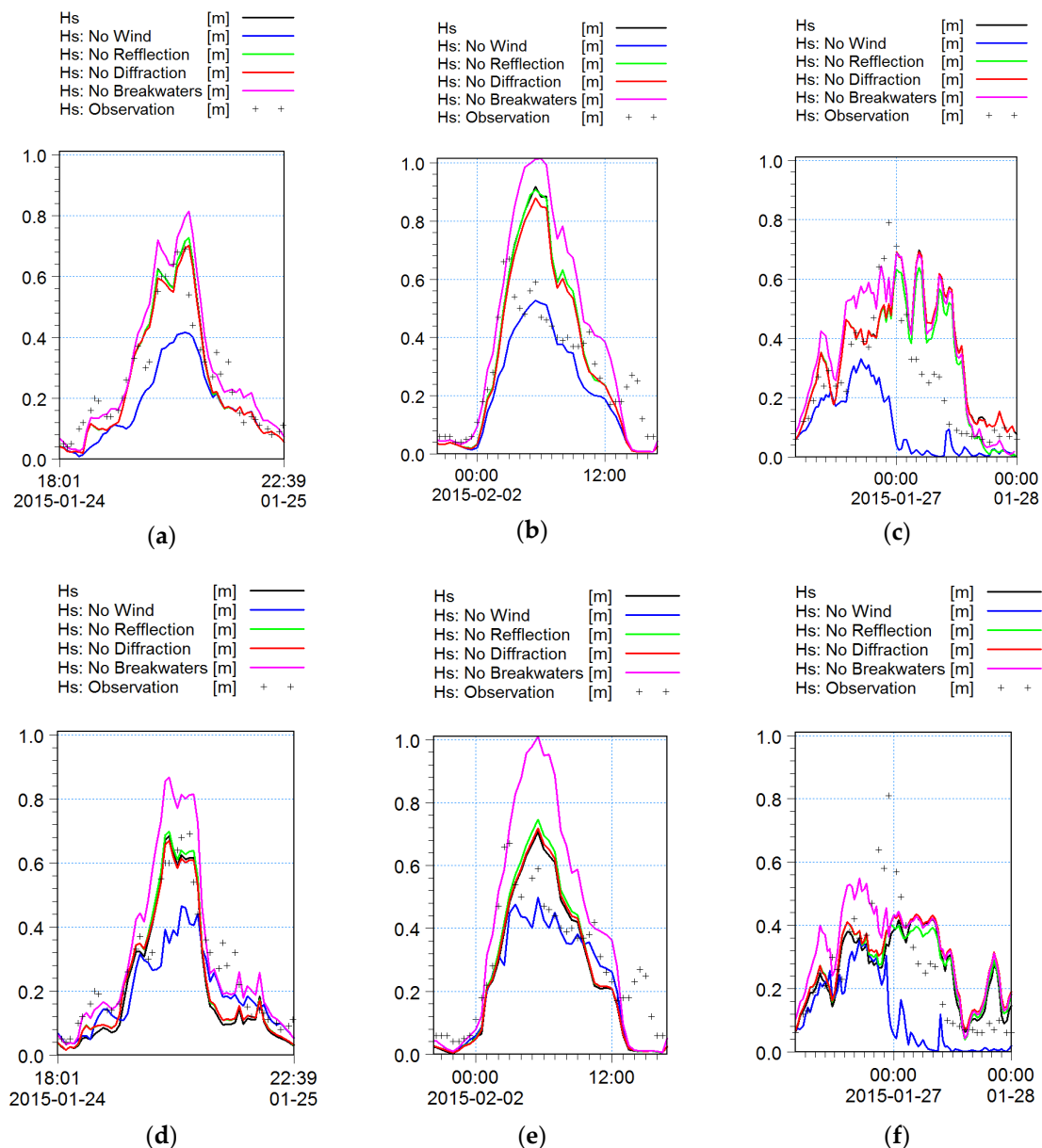


Figure 13. Sensitivity of the models to wind, reflection, diffraction module, and the breakwaters. Black: the model results considering all processes. Blue: The model results without wind forcing. Green: the model results without reflection from the breakwaters. Orange: the model results with disabling diffraction module. Pink: the model results with removing the breakwaters. (a) SWAN, 25 January; (b) SWAN, 2 February; (c) SWAN, 27 January; (d) MIKE21SW, 25 January; (e) MIKE21SW, 2 February; (f) MIKE21SW, 27 January.

In some events, such as 6 April and 6 March, a delay in the growth of the waves was observed in the models results, however, the delay in SWAN results was about one hour more than MIKE21SW results. There were three events (27 January, 5 February, and 15 February) with high winds from the north and the models did not deliver good simulations in two, 27 January and 15 February. These events were accompanied with 90-degree rotation of wind vectors during the storm. This inconsistency may be a consequence of the uncertainty in the difference between the wind observed at the buoy and that over the harbor. The wind observation station is located about 20 km away from the coastline. When wind blow from the coast (from the north), the wind boundary layer near the coastline become very complex. This complicity may not be observed at the wind station located 20 km away from the coastline. In addition, the significant wave height evaluation outside the harbor for northerly storms,

Figure 11, demonstrates that the models overestimate the significant wave height when winds were from the north. It implies that, for northern storms, winds at CLIS are more intensive than New Haven. Also, New Haven estuary has a complex geometry. Consequently, a small change in wind direction can change the fetch length drastically. For example, for NH2 station, a wind blowing from 10 degrees has more than double fetch length of the wind from the north, (Figure 1). The inconsistency of the models with observation for a very small fetch length with strong wind condition needs more investigation. This requires a local high frequency wind and wave observation to record precisely oscillation in wind speed and direction.

Another possible source of error is that a spectral wave model cannot simulate wave resonance due to reflected waves from the harbor edges and breakwaters. Figure 4b,d show spectra for 27 January and 15 February, when the models failed to accurately simulate the wave field, there were two waves, a low frequency wave from the south and high frequency waves propagating from the north. In this situation, reflection could have significant role. To observe these effects more accurately, multiple high frequency wave observation stations in the harbor would be required. The wave spectra shown in Figure 4 are not confined to a narrow range of directions and, in some events such as 2 February, the spectrum has multiple peaks. For this storm SWAN overestimated the wave height by 50%. Both models under-predicted significant wave heights during the storm on 7 March. In this event, the wave spectrum, also, have two peaks (Figure 4e). The models had good performance for the two-peak storm on 25 January, however, this storm had two peaks from the same direction but different frequency. The models may have lower accuracy in the events that the wave spectrum have multiple peaks at different directions, but that is not always the case as MIKE21SW well predicted the storm on 2 February with multiple spectrum peaks. Further investigation is required.

The influence of currents on waves was not considered in this study. Current can be considered in the modeling using a coupled hydrodynamic-wave models such as FVCOM-SWAVE, ADCIRC-SWAN, XBEACH [33,34], and MIKE21 Coupled Model. This effect should be considered in modeling in forthcoming studies. Another uncertainty in the results can be due to uncertainty in the model wave boundary condition. Though we used the observed wave spectra, instead of a model or parametrization, to prescribe the open boundary condition, it is possible that there are periods when there is variation in wave conditions along the boundary. Multiple wave sensors would be needed to assess that possibility and we hope to evaluate it in the future.

Besides the accuracy of the model, the efficiency and simplicity of the models are important in assessment of the model utility. SWAN is much more efficient computationally than MIKE21SW. SWAN executes the same model grid and configuration using the same computational engine much faster than MIKE21SW. Although SWAN is equipped with a fast-computational algorithm, the total computational time of the models in the optimum conditions was almost the same. SWAN requires much finer mesh to reach the optimum condition. In this study, the grid size of the mesh used for SWAN ($dx_{max} = 250$ m and $dx_{min} = 25$ m) was half of the grid size of the mesh used for MIKE21SW ($dx_{max} = 500$ m and $dx_{min} = 50$ m). Getting to the optimum condition with MIKE21SW was faster than SWAN. MIKE21SW automatically selects the optimum time step based on the grid size. User just needs to assign the minimum and maximum time steps. In addition, MIKE21SW showed less sensitivity to grid size than SWAN. Therefore, finding the optimum condition in MIKE21SW needs fewer number of sensitivity simulations than SWAN.

The time integration and spatial discretization method employed in SWAN and MIKE21SW play the main role in determining the differences in the results and efficiency. MIKE21SW uses an explicit Euler scheme for time integration with cell-centered finite volume when computing wave propagation while SWAN uses a fully implicit method for time integration with finite difference first order BSBT scheme. This study suggests that the method used in SWAN is computationally much faster but it is more sensitive to spatial resolution and requires much finer mesh.

The application of the spectral models for simulation of waves inside the harbor in the presence breakwater is questioned by some authors [1,2,35]. This implies that the spectral models are not suitable

for small ports with narrow connections to the ocean. However, in situations in which breakwaters have low reflection coefficients, are far enough the shoreline and do not shade a significant portion of the basin [1] spectral models may be useful as we have demonstrated in this work. In situ observations and sensitivity analyses are extremely valuable in assessing model effectiveness at all sites.

6. Conclusions

SWAN and MIKE21SW are two spectral wave models that solve the wave action balance equations. Although there are lots of similarities in both the main equations and wave source terms, they have some minor differences in the algorithms used to obtain solutions that impact both the results and efficiency of the models. SWAN and MIKE21SW were assessed on the unstructured grid and inside a harbor in the presence of three detached breakwaters. This study suggests the results of the models were consistent with observations during the storms which were affected by breakwaters. The R^2 was approximately 0.6 for both models. Considering the complexity of the modeling domain, the results are quite acceptable. The models behaved similarly in most events, MIKE21SW slightly better simulated significant wave height at storm peaks in some events. SWAN required the finer grid to get to the optimum condition, but as it uses the faster computational algorithm, the total computational time for their optimum condition was almost the same. MIKE21SW automatically selects the efficient time step based on grid size and it was less sensitive to grid size than SWAN. Therefore, the optimum condition of MIKE21SW was reached with fewer sensitivity simulations. Both models performed poorly for when high wind blew from the coast to sea. It is likely that this was a consequence of inadequate resolution of the wind field, though further observations and investigations will be required to fully understand that result. The sensitivity analysis demonstrates the wind effect was significant on the results due to large fetch length in the harbor. This is also the reason for using the spectral models for this case study. Also, it has been shown that, in MIKE21SW simulation, the breakwaters dissipate wave energy slightly more than the breakwaters in SWAN simulation.

Author Contributions: Conceptualization, A.I. and J.O.; Methodology, A.I.; Software, A.I.; Validation, A.I. and J.O.; Formal Analysis, A.I.; Investigation, A.I. and J.O.; Resources, A.I.; Writing—Original Draft Preparation, A.I.; Writing—Review & Editing, A.I. and J.O.; Visualization, A.I.; Supervision, J.O.; Project Administration, J.O.; Funding Acquisition, J.O.

Funding: The observations in this research were funded by NOAA and the Connecticut Department of Energy and Environmental Protection.

Acknowledgments: We are grateful to DHI for providing us with the MIKE21 package, and TU Delft for making SWAN open source. Connecticut Institute for Resilience and Climate Adaptation (CIRCA) supported the observations with funding from NOAA and the Connecticut Department of Energy and Environmental Protection. We are also grateful to Kay Howard-Strobel, Ale Cifuentes, and Molly James for valuable technical assistance.

Conflicts of Interest: The authors declare no conflict of interest.

References

1. Holthuijsen, L.H.; Herman, A.; Booij, N. Phase-decoupled refraction-diffraction for spectral wave models. *Coast. Eng.* **2003**, *49*, 291–305. [[CrossRef](#)]
2. Boshek, M.R. Reflection and Diffraction around Breakwaters. Master's Thesis, Delft University of Technology, Department of Hydraulic Engineering, Delft, The Netherlands, 2009.
3. Dietrich, J.C.; Tanaka, S.; Westerink, J.J.; Dawson, C.N.; Luettich, R.A., Jr.; Zijlema, M.; Holthuijsen, L.H.; Smith, J.M.; Westerink, L.G.; Westerink, H.J. Performance of the Unstructured-Mesh, SWAN+ADCIRC Model in Computing Hurricane Waves and Surge. *J. Sci. Comput.* **2012**, *52*, 468–497. [[CrossRef](#)]
4. Zijlema, M. Parallel, unstructured mesh implementation for SWAN. *Coast. Eng.* **2009**, *5*, 470–482. [[CrossRef](#)]
5. Strauss, D.; Mirferendesk, H.; Tomlinson, R. Comparison of two wave models for Gold Coast, Australia. *J. Coast. Res.* **2007**, *50*, 312–316.
6. Moeini, M.H.; Etemad-Shahidi, A. Application of two numerical models for wave hindcasting in Lake Erie. *Appl. Ocean Res.* **2007**, *29*, 137–145. [[CrossRef](#)]

7. Fonseca, R.B.; Gonçalves, M.; Guedes Soares, C. Comparing the Performance of Spectral Wave Models for Coastal Areas. *J. Coast. Res.* **2017**, *33*, 331–346. [[CrossRef](#)]
8. Hoque, M.A.; Perrie, W.; Solomon, S.M. Evaluation of two spectral wave models for wave hindcasting in the Mackenzie Delta. *Appl. Ocean Res.* **2017**, *62*, 169–180. [[CrossRef](#)]
9. Booij, N.; Ris, R.C.; Holthuijsen, L.H. A third-generation wave model for coastal regions: 1. model description and validation. *J. Geophys. Res. Oceans* **1999**, *104*, 7649–7666. [[CrossRef](#)]
10. DHI group. *MIKE 21 Spectral Wave Module. Scientific Documentation*; Danish Hydraulic Institute (DHI): Hørsholm, Denmark, 2017; 56p.
11. Mei, C.C. *The Applied Dynamics of Ocean Surface Waves*; Volume 1 of Advanced Series on Ocean Engineering; World Scientific: Singapore, 1989; ISBN 9971507897.
12. Komen, G.J.; Cavaleri, L.; Donelan, M.; Hasselmann, K.; Hasselmann, S.; Janssen, P.A.E.M. *Dynamics and Modelling of Ocean Waves*; Cambridge University Press: Cambridge, UK, 1994.
13. Young, I.R. *Wind Generated Ocean Waves, Elsevier Ocean Engineering Series*, 1st ed.; Bhattacharyya, R., McCormick, M.E., Eds.; Elsevier: Amsterdam, The Netherlands, 1999; Volume 2, ISBN 9780080433172.
14. SWAN team. *Swan Scientific and Technical Documentation*; Delft University of Technology: Delft, The Netherlands, 2018; 147p.
15. Janssen, P.A.E.M. Wave-Induced Stress and the Drag of Air Flow over Sea Waves. *J. Phys. Oceanogr.* **1989**, *19*, 745–754. [[CrossRef](#)]
16. Janssen, P.A.E.M. Quasi-linear Theory of Wind-Wave Generation Applied to Wave Forecasting. *J. Phys. Oceanogr.* **1991**, *21*, 1631–1642. [[CrossRef](#)]
17. Gunther, H.; Hasselmann, S.; Janssen, P.A.E.M. *The WAM Model Cycle 4 (Revised Version)*; Deutsch. Klim. Rechenzentrum, Techn. Rep. No. 4; Deutsches Klimarechenzentrum: Hamburg, Germany, 1992.
18. Komen, G.J.; Hasselmann, K.; Hasselmann, K. On the Existence of a Fully Developed Wind-Sea Spectrum. *J. Phys. Oceanogr.* **1984**, *14*, 1271–1285. [[CrossRef](#)]
19. WAMDI Group. The WAM Model—A Third Generation Ocean Wave Prediction Model. *J. Phys. Oceanogr.* **1988**, *18*, 1775–1810. [[CrossRef](#)]
20. Yan, L. *An Improved Wind Input Source Term for Third Generation Ocean Wave Modelling*; Scientific report WR-No 87-8; Royal Netherlands Meteorological Inst.: De Bilt, The Netherlands, 1984.
21. Hasselmann, K. On the spectral dissipation of ocean waves due to white capping. *Bound. Layer Meteorol.* **1974**, *6*, 107–127. [[CrossRef](#)]
22. Van der Westhuysen, A.J. Advances in the Spectral Modelling of Wind Waves in the Nearshore. Ph.D. Thesis, Delft University of Technology, Department of Civil Engineering, Delft, The Netherlands, 2007.
23. Van der Westhuysen, A.J.; Zijlema, M.; Battjes, J.A. Nonlinear saturation-based whitecapping dissipation in SWAN for deep and shallow water. *Coast. Eng.* **2007**, *54*, 151–170. [[CrossRef](#)]
24. Hasselmann, K.; Barnett, T.P.; Bouws, E.; Carlson, H.; Cartwright, D.E.; Enke, K.; Ewing, J.A.; Gienapp, H.; Hasselmann, D.E.; Kruseman, P.; et al. *Measurements of Wind-Wave Growth and Swell Decay during the Joint North Sea Wave Project (JONSWAP)*; Ergänzungsheft 8–12; Deutsches Hydrographisches Institut: Hamburg, Germany, 1973.
25. Zijlema, M.; van Vledder, G.P.; Holthuijsen, L.H. Bottom friction and wind drag for wave models. *Coast. Eng.* **2012**, *65*, 19–26. [[CrossRef](#)]
26. Battjes, J.A.; Janssen, J.P.F.M. Energy loss and set-up due to breaking of random waves. In Proceedings of the 16th International Conference Coastal Engineering, Hamburg, Germany, 27 August–3 September 1978; pp. 569–587. [[CrossRef](#)]
27. Eldeberky, Y.; Battjes, J.A. Parameterization of triad interactions in wave energy models. In Proceedings of the Coastal Dynamics Conference '95, Gdansk, Poland, 4–8 September 1995; pp. 140–148.
28. Eldeberky, Y. Nonlinear Transformation of Wave Spectra in the Nearshore Zone. Ph.D. Thesis, Delft University of Technology, Department of Civil Engineering, Delft, The Netherlands, 1996.
29. Hasselmann, S.; Hasselmann, K.; Allender, J.H.; Barnett, T.P. Computations and Parameterizations of the Nonlinear Energy Transfer in a Gravity-Wave Spectrum. Part II: Parameterizations of the Nonlinear Energy Transfer for Application in Wave Models. *J. Phys. Oceanogr.* **1985**, *15*, 1378–1391. [[CrossRef](#)]
30. U.S. Army Corps of Engineers. *Coastal Engineering Manual Part VI*; Books Express Publishing: Newbury, UK, 2012.

31. Seelig, W.N. Wave Reflection from Coastal Structures. In Proceedings of the Coastal Structures '83, Arlington, Virginia, 9–11 March 1983; American Society of Civil Engineers: Reston, VA, USA; pp. 961–973.
32. Davidson, M.A.; Bird, P.A.; Bullock, G.N.; Huntley, D.A. Wave Reflection: Field Measurements, Analysis and Theoretical Developments. In Proceedings of the Coastal Dynamics '94, Barcelona, Spain, 21–25 February 1994; American Society of Civil Engineers: Reston, VA, USA; pp. 642–655.
33. Sánchez-Arcilla, A.; García-León, M.; Gracia, V. Hydro morphodynamic modelling in Mediterranean storms: Errors and uncertainties under sharp gradients. *Nat. Hazards Earth Syst. Sci.* **2014**, *14*, 2993–3004. [[CrossRef](#)]
34. Gracia, V.; García-León, M.; Grifoll, M.; Sánchez-Arcilla, A. Breaching of a barrier under extreme events: The role of morphodynamic simulations. *J. Coast. Res.* **2013**, *65*, 951–956. [[CrossRef](#)]
35. Booij, N.; Holthuijsen, L.H.; de Lange, P.H.M. The penetration of short-crested waves through a gap. In Proceedings of the 23rd International Conference on Coastal Engineering, Venice, Italy, 4–9 October 1992; pp. 1044–1052.



© 2018 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>).